

Conservative Fields: Gravity, Electrostatic, and Magnetic fields follow the **Poisson's equation** of the form:

 $\nabla^2 \psi = f$  (sources)

where  $\psi$  is a **potential**,  $\vec{\nabla}$  is the gradient operator:

$$\vec{\nabla} = \frac{\partial}{\partial x}\hat{i} + \frac{\partial}{\partial y}\hat{j} + \frac{\partial}{\partial z}\hat{k}$$

$$\nabla^2 = \vec{\nabla} \cdot \vec{\nabla} = \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2}$$

**Notation**: Here, the arrow → denotes a vector quantity; the carat ^ denotes a unit direction vector.

Hence, the gradient operator is just a vector form of slope...

Because **Poisson's** eqn. always incorporates a potential  $\psi$ , we call these **Potential Field Methods** 

### **Potential Field Theory**

For any CONSERVATIVE field,  $F: \vec{F} = \vec{\nabla} \psi$  (1)

Fluxes: Johann Carl Friedrich Gauss (1777 – 1855):  
Gauss ' Law for gravity: 
$$\Phi_g = \bigoplus_{s} \vec{g} \cdot \vec{n} \, dA = -4 \pi GM$$
 (2 a)  
Gauss ' Law for electrostatics:  $\Phi_E = \bigoplus_{s} \vec{E} \cdot \vec{n} \, dA = \frac{q}{\epsilon_0}$  (2 b)  
Gauss ' Law for magnetism:  $\Phi_B = \bigoplus_{c} \vec{B} \cdot \vec{n} \, dA = 0$  (2 c)

But by Gauss ' Divergence Theorem (for gravity):  

$$\Phi_{g} = \bigoplus_{s} \vec{g} \cdot \vec{n} \, dA = \bigoplus_{v} \vec{\nabla} \cdot \vec{g} \, dV \qquad (3)$$
Also,
$$M = \bigoplus_{v} \rho \, dV \qquad (4)$$

### **Potential Field Theory**

Thus, (2a), (3), (4) yield:  

$$\Phi_g = \bigoplus_V \vec{\nabla} \cdot \vec{g} \, dV = -4\pi G M = -4\pi G \bigoplus_V \rho \, dV \qquad (5)$$

Therefore, 
$$\vec{\nabla} \cdot \vec{g} = -4\pi G \rho$$
 (6)

But (1) 
$$\Rightarrow \quad \vec{g} = \vec{\nabla}\psi$$
, yielding: (7)

$$\vec{\nabla} \cdot (\vec{\nabla} \psi) = -4\pi G \rho = \Im \text{ (sources)}$$
(8)

that is,

$$\nabla^2 \psi = -4\pi G \rho = \Im$$
, the POISSON's equation (9)

which is similar for the electrostatic / magnetic fields.

Applying (2a) for a spherical mass distribution of uniform density  $\rho$ , radius, r:

$$\oint_{S} \vec{g} \cdot \vec{n} \, dA = g(4\pi r^{2}) = -4\pi GM$$

$$\Rightarrow \quad g = -\frac{GM}{r^{2}} \quad (Newton's \, Law!)$$
(10)

## Gravity

**Gravitational field**  $\vec{g}$ :

$$\vec{\nabla}\Psi = \vec{g}$$

Poisson's equation,

$$\nabla^2 \Psi = \vec{\nabla} \cdot \vec{g} = -4\pi G \rho$$

given a mass density of  $\rho$ , where G is universal gravitational constant:

$$G = 6.672 \times 10^{-11} \text{ Nm}^2/\text{kg}^2$$

Gravitational Flux,  $\Phi$ :

**Flux,** 
$$\Phi = -4 \pi G \iiint_{Vol} \rho dV$$

#### **Newton's Law of Gravitation**

Thus, for a body with mass *M* is spherical with constant density:

$$\vec{g} = \hat{r} \frac{GM}{r^2}$$

Where, r is distance from the center of mass, and  $\hat{r}$  is the (unit) direction vector pointing toward the center.

Kepler's 3<sup>rd</sup> Law:  $T^2 = kr^3$  (1619)  $\rightarrow$  Newton's Law of Gravitation (1687): Kepler's constant,  $k \approx (\sqrt[3]{4}) \times 10^{-5} (day)^2 / (A.U.)^3 \approx 1 (yr)^2 / (A.U.)^3$  !!! From Newton's Law, equating Centripetal & gravitational forces,  $k = 4\pi^2 / (GM)$ 

Now, 
$$\vec{F} = \hat{r} \frac{GMm}{r^2} = m\vec{g}$$

So  $\vec{g}$  expresses the acceleration of *m* due to *M*, and has units of acceleration: **Gal** in cgs (= 0.01 m/s<sup>2</sup> ≈ 1 milli-g)

On the Earth's surface,  

$$\vec{g} = \hat{r} \frac{GM_E}{R_E^2} = \hat{r} \frac{(6.67 \times 10^{-11})(5.97 \times 10^{24})}{(6.38 \times 10^6)^2} = 9.8 \ m/s^2$$

**HOWEVER**,  $\rho$  is **not** radially symmetric in the Earth... so  $\vec{g}$  is neither constant or uniform!

Gravity methods look for *anomalies*, or perturbations, from a reference value of  $\vec{g}$  at the Earth's surface:



## **Example:**

# Global Free-Air Gravity Field from GRACE + GOCE + satellite altimetry + surface measurements...

Surface Free-Air Anomaly



WGM-2012 model from Bureau Gravimetríque International

### **Gravity Anomalies from GRACE**

